# Assessment of steady-state seepage through dams with nonsymmetric boundary conditions: analytical approach 

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#### Abstract

In this study, new analytical solutions were developed for 2D and 3D steady-state water seepage through dams with nonsymmetric boundary conditions. The nonsymmetric boundary conditions for the 2D cases were created with different unit step functions on a part and/or parts of the right boundary of dam plane. Six cases were investigated in 2D, where a constant hydraulic head is applied at the left boundary of the dam plane and rectangular, ramp, triangular, trapezoidal, tunnel, and piecewise rectangular distributions of hydraulic head are applied at the right boundary of the dam plane. Then, a 3D case with a constant hydraulic head at the upstream and a linearly distributed hydraulic head at the downstream of the dam was investigated. Subsequently, the performance of proposed analytical solutions was examined by comparison with numerical finite difference modeling. The results demonstrate reasonable accuracy of the developed equations. The developed analytical solutions can be utilized as a benchmark to verify numerical models with similar boundary conditions.


Keywords Dam • Seepage • Analytical approach . Nonsymmetric boundary condition • Partial differential equation

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## Introduction

Water seepage through the dam body has a vital role in the dam stability analysis. Therefore, hydraulic head determination and subsequently seepage flow simulation are among the design requisites in these structures. Water seepage into earth dams is illustrated by unconfined aquifer equation, a nonlinear partial differential equation (PDE) with limited analytical solutions (Zarif Sanayei et al. 2015, 2019). However, the choice of steady-state flow and homogeneous isotropic soil for unconfined aquifer make it feasible to acquire analytical solutions through classical approaches.

Tan et al. (2017) studied the effect of spatial variability of hydraulic parameters on the seepage flow in earth dams (Javdanian et al. 2018 a, b: Javdanian and Pradhan 2019; Shakarami et al. 2019; Javdanian 2019; Nasiri et al. 2019). Seepage flow through dams have been investigated by many researchers through finite difference (e.g., Fu and Jin 2009; Rakhshandehroo and Pourtouiserkani 2013), boundary elements (e.g., Chen et al. 1994; Leontiev and Huacasi 2001), and mesh-free numerical (e.g., Navas et al. 2016; Zhang et al. 2017) methods. Li et al. (2017) analyzed the seepage in roller compacted concrete dam using the finite element method (FEM). Nourani et al. (2014) proposed an arrangement of tubes analog and multiple reservoirs to simulate seepage flow through embankment dam conducting laboratory studies. An efficient technique to capture the seepage face boundary condition in finite element modeling of seepage problems in porous media (Fusi et al. 2015) was proposed by Pedroso (2015). Yuan and

Zhong (2016) developed a quadrature element formulation for 3D analysis of unconfined seepage in earth dams. A three-dimensional (3D) numerical analysis for unconfined seepage problems in anisotropic and inhomogeneous domains was developed by KazemzadehParsi and Daneshmand (2013). Fukuchi (2016) used the interpolation finite difference method (FDM) for numerical modeling of steady-state seepage problems. They reported that the proposed method is effective in the flow net calculation under mixed Dirichlet and Neumann conditions.

Analytical solutions give a better insight in comparison to the numerical modeling (Kacimov and Obnosov 2012, 2019; Wei et al. 2018). Kacimov and Obnosov (2012) developed analytical solutions for 2D seepage in dam core with uniform hydraulic head. Rezk and Senoon (2012) presented an analytical solution to calculate seepage through earth dam with upstream blanket. Teloglou and Bansal (2012) presented analytical solutions of Boussinesq equation to predict flow rate of water in an unconfined aquifer under transient seepage. Liang and Zhang (2013) derived analytical solutions for the lateral discharge and water table in a heterogeneous unconfined aquifer with fluctuating river stage. Closedform analytical solutions to capture the influence of sea level rising on the unconfined sloping island aquifers were developed by Chesnaux (2016). El Tani et al. (2019) developed analytical equations for water table drawdown. Some analytical methods were also proposed by Mohsenian et al. (2019) for two-dimensional modeling of groundwater flow.

Reviewing the available literature indicates that most of the studies in the field of water seepage have been performed on the problems with symmetric boundary conditions by numerical methods. The current research focuses on the analytical solutions of two- and three-dimensional steady-state water seepage through dams with nonsymmetric boundary conditions. The boundary conditions for 2D cases are created with different unit step functions on a part and/or parts of the right boundary of the dam plane. In 2D cases, a constant hydraulic head is applied at the left boundary and various hydraulic head distributions are applied at the right boundary of the dam plane. Then, a 3D case is investigated. Subsequently, the performances of developed analytical solutions are examined using numerical modeling.

## Governing equations

The equation for water flow in unconfined aquifer is developed by combining continuity and Darcy's law as a momentum equation. The equation for Cartesian coordinate in the inhomogeneous anisotropic soil with Dupuit assumptions is (Serrano and Workman 1998; Fukuchi 2016):
$\frac{\partial h}{\partial t}=\frac{1}{S_{y}} \frac{\partial}{\partial x}\left(K_{x} h \frac{\partial h}{\partial x}\right)+\frac{1}{S_{y}} \frac{\partial}{\partial y}\left(K_{y} h \frac{\partial h}{\partial y}\right)$
where $h(L)$ is the hydraulic head, $S_{y}$ is the specific yield, $t$ $(T)$ is the time, and $K_{x}$ and $K_{y}\left(\frac{L}{T}\right)$ are the hydraulic conductivity of soil in $x$ and $y$ directions, respectively.

For an unconfined aquifer with steady-state flow in the isotropic homogeneous soil, Eq. (1) is written as:
$\frac{\partial^{2} h^{2}}{\partial x^{2}}+\frac{\partial^{2} h^{2}}{\partial y^{2}}=0$
Analytical solutions of Eq. (2) would express spatial variability of the hydraulic head in the unconfined aquifer sample under the given boundary conditions. In this research, new two-dimensional (2D) analytical solutions are developed for Eq. (2) subjected to nonsymmetric boundary conditions.

## Analytical solution for 2D water seepage

In order to obtain closed-form analytical solution for Eq. 2, some simplifications need to be employed. A schematic rectangular cube of a homogeneous dam with various boundary and initial conditions is considered (Fig. 1a). Then, closed-form analytical solutions are investigated for the 2D horizontal plane (Fig. 1b) with nonsymmetric hydraulic head condition on the boundaries. In particular, constant hydraulic head is often applied on the left boundary and various hydraulic head distributions are applied on the right boundary (Fig. 1a, b).

## Hydraulic head distribution

## Rectangular distribution

The case is applicable where the outlet of a dam passes through a rectangular conduit and discharges into a channel, which creates a constant head at a part of the downstream of the dam. A constant hydraulic head is applied at the left boundary and a rectangular distribution of

Fig. 1 a Schematic rectangular cube of a homogeneous dam. b 2 D horizontal plane of the dam

hydraulic head is applied at the right boundary of the dam plane. In addition, the top and the bottom boundaries have no-flow boundary conditions (Fig. 1b). The boundary conditions for are expressed as:
$h(0, y)=h_{1}$
$h(L, y)=\frac{h_{1}}{a} *\left[U(y)-U\left(y-\frac{B}{10}\right)\right]$
$\frac{\partial h}{\partial y}(x, 0)=0 \quad \frac{\partial h}{\partial y}(x, B)=0 \quad(3-\mathrm{c})$
where $U$ is the unit step function, and $h_{1}$ and $a$ are constants.

A schematic view of the right boundary condition (Eq. 3-b) with $a=3, h_{l}=300 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ is depicted in the Fig. 2. As depicted in this figure, the boundary condition presents the constant hydraulic head for all points at $0 \leq y \leq \frac{B}{10}$ and creates zero hydraulic head in other points. To find an analytical solution for Eq. (2) with the aforementioned boundary conditions, the variable change $h^{2}(x, y)=w(x, y)$ is used. Then, the partial differential equation (PDE) and the boundary conditions for $w(x, y)$ are written as:
$\frac{\partial^{2} w}{\partial x^{2}}+\frac{\partial^{2} w}{\partial y^{2}}=0$

Fig. 2 Right boundary condition (Eq. 3-b) for the rectangular distribution case

$w(0, y)=h_{1}{ }^{2}$
$w(L, y)=\left(\frac{h_{1}}{a}\right)^{2} *\left[U(y)-U\left(y-\frac{B}{10}\right)\right]^{2}$
$(4-\mathrm{b}) \quad$ tions are obtained for $\frac{Y^{\prime \prime}}{Y}$ as follows:
$(4-\mathrm{c}) \quad-\frac{X^{\prime \prime}}{X}=0$
$\frac{\partial w}{\partial y}(x, 0)=0 \quad \frac{\partial w}{\partial y}(x, B)=0$
$-\frac{X^{\prime \prime}}{X}=-\lambda^{2}$
Using separation of variables for $w(x, y)$ :
$w(x, y)=X(x) Y(y)$
Then, substituting Eq. (5) into Eq. (4-a) and employing some simplifications gives:
$\frac{Y^{\prime \prime}}{Y}=-\frac{X^{\prime \prime}}{X}=\mu$
where $\mu$ is an arbitrary constant.
If $\mu=0$ and $\mu<0$, assuming $\mu=-\lambda^{2}$ and $\lambda>0$, then, $Y(y)$ in Eq. (6) has two answers considering the boundary conditions of Eq. (4-d), as follows:
$Y_{1}=B_{1}$
$Y_{n}=C_{1} \cos (\lambda y), \quad \lambda=\frac{n \pi}{B} \quad n=1,2,3, \ldots, \infty$
where $B_{1}$ and $C_{1}$ are constants.
where $C_{2}^{*}$ is $B_{1} C_{2}$, and $C_{3}^{*}$ is $B_{1} C_{3}$.
Also, by substituting Eq. (10-b) and Eq. (8) into Eq. (5), the second answer for $w(x, y)$ is written as:

$$
\begin{align*}
w_{2}(x, y) & =X_{n}(x) Y_{n}(y) \\
& =\sum_{n=1}^{\infty} \cos (\lambda y)\left[C_{4}^{*} \sinh (\lambda x)+C_{5}^{*} \cosh (\lambda x)\right] \tag{12}
\end{align*}
$$

where $C_{4}^{*}$ is $C_{1} C_{4}$, and $C_{5}^{*}$ is $C_{1} C_{5}$.
Finally, the overall answer for $w(x, y)$ is the combination of Eq. (11) and Eq. (12):

$$
\begin{aligned}
w(x, y) & =w_{1}(x, y)+w_{2}(x, y) \\
& =C_{2}^{*} x+C_{3}^{*}+\sum_{n=1}^{\infty} \cos (\lambda y)\left[C_{4}^{*} \sinh (\lambda x)+C_{5}^{*} \cosh (\lambda x)\right]
\end{aligned}
$$

By substituting the boundary conditions of Eq. (4-b) and Eq. (4-c) into Eq. (13) and using Fourier series properties for Eq. (13), the constant coefficients in Eq. (13) are written as:

$$
\begin{align*}
& C_{2}^{*}=\frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left[U(y)-U\left(y-\frac{B}{10}\right)\right]^{2} \quad(14-\mathrm{a})  \tag{14-a}\\
& d y-\frac{C_{3}^{*}}{L}=\frac{\frac{1}{10}\left(\frac{h_{1}}{a}\right)^{2}-h_{1}^{2}}{L} \\
& C_{3}^{*}=\frac{1}{B} \int_{0}^{B} h_{1}^{2} d y=h_{1}^{2}  \tag{14-b}\\
& C_{4}^{*}=\frac{\frac{2}{B} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} *\left[U(y)-U\left(y-\frac{B}{10}\right)\right]^{2} \cos (\lambda y) d y-C_{5}^{*} \cosh (\lambda L)}{\sinh (\lambda L)}(14-\mathrm{c}) \\
&=\frac{2 h_{1}^{2} \sin \left(\frac{1}{10} n \pi\right)}{a^{2} n \pi \sinh \left(\frac{n \pi L}{B}\right)} \\
& C_{5}^{*}=\frac{2}{B} \int_{0}^{B} h_{1}^{2} \cos (\lambda y) d y=0 \tag{14-d}
\end{align*}
$$

Substituting Eq. (13) into $h^{2}(x, y)=w(x, y)$ gives:
$h(x, y)=\sqrt{C_{2}^{*} x+C_{3}^{*}+\sum_{n=1}^{\infty} \cos (\lambda y)\left[C_{4}^{*} \sinh (\lambda x)+C_{5}^{*} \cosh (\lambda x)\right]}$

Exact analytical solution of the problem (Eq. 15) satisfies the PDE (Eq. 2) and the boundary conditions (Eq. 3a-c). To confirm convergence of summations in Eq. (15), hydraulic head at different locations is computed using summations truncation with different amounts of $n$ (Table 1). This table is presented for the parameters $a=3, h_{l}=300 \mathrm{~cm}, L=100 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$. As shown in Table 1, the change in the hydraulic head is
negligible at different values of $n$. Also, the values are approximately equal at $n=1 \sim 15$ and $n=1 \sim 20$. To illustrate the application of the derived equations, hydraulic head values from an explicit scheme FDM solution (Eq. 2) is compared to the analytical solution (Eq. 15) for various amounts of $x$ and $y$ (Table 1, columns 7). Relative error ( $R E$ ) (Javdanian and Lee 2019; Javdanian 2019) in Table 1 is calculated using Eq. (16):
$R E=\frac{\left|h_{\text {analytical }}-h_{\text {numerical }}\right|}{\left|h_{\text {analytical }}\right|} \times 100 \%$
where, $h_{\text {analytical }}$ and $h_{\text {numerical }}$ are analytical-based and numerical-based hydraulic head, respectively.

As shown in Table 1, the values of $R E$ are less than $3 \%$. Figure 3 depicted hydraulic head contours calculated by Eq. (15). This figure shows the distribution of the hydraulic head that remains at $h=300 \mathrm{~cm}$ on the left $(x=0)$ and $h=$ Eq. (3-b) on the right boundaries ( $x=$ 100 cm ) of the dam plane. Hydraulic head contours are perpendicular to the top and bottom boundaries of the dam plane, confirming no-flow boundary condition.

## Ramp distribution

Ramp distribution of the hydraulic head on the downstream of the dam body may be caused by the right triangular shape of the water outlet. In practice, this condition occurs during clogging of the sluice gate by sediments and debris. Mathematically, top and bottom, as well as the left boundary conditions, are as same as the rectangular distribution case. The right boundary is written as Eq. (17):

$$
\begin{align*}
h(L, y)= & \frac{h_{1}}{a} \times \frac{3}{B} \times\left(y-\frac{B}{3}\right) \times U\left(y-\frac{B}{3}\right) \\
& \times\left[U\left(y-\frac{B}{3}\right)-U\left(y-\frac{2 B}{3}\right)\right] \tag{17}
\end{align*}
$$

The right boundary condition (Eq. 17) with $a=3$, $h_{l}=300 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ was illustrated in Fig. 4. The boundary condition presents the ramp distribution of hydraulic head for $\frac{B}{3} \leq y \leq \frac{2 B}{3}$. Following similar mathematical methods used for rectangular distribution case, the answer for $h(x, y)$ would be identical to Eq. (15) but with different coefficients. Mathematical forms of $C_{3}^{*}$ and $C_{5}^{*}$ are identical to the ones defined by Eq. (14); however, $C_{2}^{*}$ and $C_{4}^{*}$ are expressed as:

Table 1 Hydraulic head values resulted from analytical approach (for different summation truncations) and FDM for rectangular distribution case

| $x(c m)$ | $y(c m)$ | $h_{\text {analytical }}$ |  | $h_{F D M}$ | $R E(\%)$ |  |  |
| :--- | :---: | :--- | :--- | :--- | :--- | :--- | :--- |
|  |  | $n=1 \sim 5$ | $n=1 \sim 10$ | $n=1 \sim 15$ | $n=1 \sim 20$ | $\Delta x=1 \mathrm{~cm}$ and $\Delta y=10 \mathrm{~cm}$ |  |
| 60 | 100 | 203.12 | 203.26 | 203.34 | 203.34 | 198.45 | 2.40 |
| 50 | 50 | 228.47 | 228.63 | 228.81 | 228.81 | 223.36 | 2.38 |
| 70 | 300 | 175.65 | 175.89 | 176.15 | 176.15 | 172.41 | 2.12 |
| 40 | 300 | 242.06 | 242.21 | 242.35 | 242.35 | 238.62 | 1.53 |
| 80 | 80 | 148.12 | 148.33 | 148.48 | 148.48 | 141.51 | 2.25 |
| 80 | 250 | 144.38 | 144.52 | 144.65 | 144.65 | 125.26 | 2.17 |
| 90 | 50 | 128.54 | 128.73 | 128.91 | 128.90 | 9.92 | 2.82 |
| 90 | 350 | 102.17 | 102.39 | 102.66 | 102.66 | 270.03 | 2.66 |
| 20 | 150 | 273.37 | 273.58 | 273.81 | 273.81 | 285.30 | 1.38 |
| 10 | 400 | 286.81 | 287.03 | 287.29 | 287.29 | 0.69 |  |

$$
\begin{align*}
C_{2}^{*}= & \frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(\frac{3}{B}\right)^{2}  \tag{18-a}\\
& \times\left[\left(y-\frac{B}{3}\right) \times U\left(y-\frac{B}{3}\right) \times\left[U\left(y-\frac{B}{3}\right)-U\left(y-\frac{2 \mathrm{~B}}{3}\right)\right]\right]^{2} \\
& d y-\frac{C_{3}^{*}}{L}=\frac{\frac{1}{9}\left(\frac{h_{1}}{a}\right)^{2}-h_{1}^{2}}{L}
\end{align*}
$$

$$
\begin{align*}
C_{4}^{*} & =\frac{\frac{2}{B} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(\frac{3}{B}\right)^{2} \times\left[\left(y-\frac{B}{3}\right) \times U\left(y-\frac{B}{3}\right) \times\left[U\left(y-\frac{B}{3}\right)-U\left(y^{-} \frac{2 B}{3}\right)\right]\right]^{2} \cos (\lambda y) d y}{\sinh (\lambda L)} \\
= & \frac{2 h_{1}^{2}\left(18 \sin \left(\frac{n \pi}{3}\right)+n^{2} \pi^{2} \sin \left(\frac{2 n \pi}{3}\right)-18 \sin \left(\frac{2 n \pi}{3}\right)+6 n \pi \cos \left(\frac{2 n \pi}{3}\right)\right)}{a^{2} n^{3} \pi^{3} \sinh \left(\frac{n \pi L}{B}\right)} \tag{18-b}
\end{align*}
$$

where, $\lambda$ is identical to the one defined for the rectangular distribution case.

Figure 5 shows the hydraulic head calculated by Eqs. (15) and (18) for $n=1 \sim 20$. As seen in this figure, the ramp distribution of the hydraulic head on the right boundary of the dam plane creates a skewed shape in the hydraulic head contours. Figure 5 depicts distribution of the hydraulic head that remains at $h=300 \mathrm{~cm}$ on the left
( $x=0$ ) and $h=$ Eq. (17) on the right boundaries ( $x=$ 100 cm ) of the dam plane.

## Triangular distribution

Triangular distribution of the hydraulic head may be resulted by the equilateral triangular shape of the water outlet. In practice, this condition occurs

Fig. 3 Hydraulic head contours for the rectangular distribution case

during clogging of the sluice gate by sediments and debris. Mathematically, bottom and top boundary conditions as well as the left boundary condition
are identical to rectangular distribution case. The right boundary is written as:
$h(L, y)=\frac{h_{1}}{a} \times\left(U\left(y-\frac{B}{5}\right) \times\left(\frac{5}{B} *\left(y-\frac{B}{5}\right)\right)+U\left(y-\frac{2 B}{5}\right) \times\left(-\frac{5}{B}\right) \times\left(\left(y-\frac{3 B}{5}\right)+\left(y-\frac{B}{5}\right)\right)+U\left(y-\frac{3 B}{5}\right)\left(\frac{5}{B} \times\left(y-\frac{3 B}{5}\right)\right)\right)$

Fig. 4 Right boundary condition (Eq. 17) for the ramp distribution case


Fig. 5 Hydraulic head contours for the ramp distribution case


The right boundary condition (Eq. 19) with $a=3$, $h_{l}=300 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ is shown in Fig. 6. As seen in Fig. 6, this boundary condition presents the triangular distribution of the hydraulic head for $\frac{B}{5} \leq y \leq$ $\frac{3 B}{5}$. The answer of $h(x, y)$ for triangular distribution case
would be identical to Eq. (15) but with different coefficients. Mathematical forms of $C_{3}^{*}$ and $C_{5}^{*}$ are identical to Eq. (14); however, $C_{2}^{*}$ and $C_{4}^{*}$ are expressed as:

$$
\begin{gather*}
C_{2}^{*}=\frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(U\left(y-\frac{B}{5}\right) \times\left(\frac{5}{B} *\left(y-\frac{B}{5}\right)\right)+U\left(y-\frac{2 B}{5}\right) \times\left(-\frac{5}{B}\right) \times\left(\left(y-\frac{3 B}{5}\right)+\left(y-\frac{B}{5}\right)\right)\right.  \tag{20-a}\\
\left.+U\left(y-\frac{3 B}{5}\right)\left(\frac{5}{B} \times\left(y-\frac{3 B}{5}\right)\right)\right)^{2} d y-\frac{C_{3}^{*}}{L}=\frac{\frac{2}{15}\left(\frac{h_{1}}{a}\right)^{2}-h_{1}^{2}}{L}
\end{gather*}
$$

$$
\begin{align*}
C_{4}^{*} & =\frac{\frac{2}{B} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} *\left(U\left(y-\frac{B}{5}\right) \times\left(\frac{5}{B} *\left(y-\frac{B}{5}\right)\right)+U\left(y-\frac{2 B}{5}\right) \times\left(-\frac{5}{B}\right) \times\left(\left(y-\frac{3 B}{5}\right)+\left(y-\frac{B}{5}\right)\right)+U\left(y-\frac{3 B}{5}\right)\left(\frac{5}{B} \times\left(y-\frac{3 B}{5}\right)\right)\right)^{2} \cos (\lambda y) d y}{\sinh (\lambda L)}  \tag{20-b}\\
& =\frac{2 h_{1}^{2}\left(50 \sin \left(\frac{n \pi}{5}\right)+20 n \pi \cos \left(\frac{2 n \pi}{5}\right)-50 \sin \left(\frac{3 n \pi}{5}\right)\right)}{a^{2} n^{3} \pi^{3} \sinh \left(\frac{n \pi L}{B}\right)}
\end{align*}
$$

The hydraulic head ( $h$ ) is calculated for $n=1 \sim 20$ using Eqs. (15) and (20) and the contours was depicted in Fig. 7. As seen in this figure, the triangular distribution of the hydraulic head on the right boundary creates a tent shape in the hydraulic head contours.

## Trapezoidal distribution

This case is applicable where the water outlet from the
dam takes place through a trapezoidal gate valves. The
This case is applicable where the water outlet from the
dam takes place through a trapezoidal gate valves. The right boundary is expressed as:

Fig. 6 Right boundary condition (Eq. 19) for the triangular distribution case


$$
\begin{align*}
& h(L, y) \\
& =\frac{h_{1}}{a} \times\left(U\left(y-\frac{B}{4}\right) \times\left(\frac{8}{B} \times\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{3 B}{8}\right) \times\left(1-\frac{8}{B}\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{5 B}{8}\right)\left(-\frac{8}{B}\left(y-\frac{6 B}{8}\right)-1\right)+U\left(y-\frac{6 B}{8}\right)\left(\frac{8}{B} \times\left(y-\frac{6 B}{8}\right)\right)\right) \tag{21}
\end{align*}
$$

The right boundary condition of the trapezoidal distribution case (Eq. 21) with $a=3, h_{l}=300 \mathrm{~cm}$, and $B=$ 500 cm was shown in Fig. 8. As depicted in this figure, the boundary condition presents the trapezoidal distribution
of the hydraulic head for $\frac{B}{4} \leq y \leq \frac{6 B}{8}$. For the trapezoidal distribution case, the mathematical forms of $C_{3}^{*}$ and $C_{5}^{*}$ are identical to Eq. (14); however, $C_{2}^{*}$ and $C_{4}^{*}$ are expressed as:

Fig. 7 Hydraulic head contours for the triangular distribution case


Fig. 8 Right boundary condition (Eq. 21) for the trapezoidal distribution case

$C_{2}^{*}=\frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2}$ $(22-a)$

$$
\begin{aligned}
& \times\left(U\left(y-\frac{B}{4}\right) \times\left(\frac{8}{B} \times\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{3 B}{8}\right) \times\left(1-\frac{8}{B}\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{5 B}{8}\right)\left(-\frac{8}{B}\left(y-\frac{6 B}{8}\right)-1\right)+U\left(y-\frac{6 B}{8}\right)\left(\frac{8}{B} \times\left(y-\frac{6 \mathrm{~B}}{8}\right)\right)\right)^{2} \\
& d y-\frac{C_{3}^{*}}{L}=\frac{\frac{1}{3}\left(\frac{h_{1}}{a}\right)^{2}-h_{1}^{2}}{L}
\end{aligned}
$$

Fig. 9 Hydraulic head contours for the trapezoidal distribution case

$C_{4}^{*}$
$(22-b)$
$=\frac{\left.\frac{2_{B}^{B}}{B} \int_{0}^{\left(\frac{h_{1}}{a}\right.}\right)^{2} *\left(U\left(y-\frac{B}{4}\right) \times\left(\frac{8}{B} *\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{3 B}{8}\right) \times\left(1-\frac{8}{B}\left(y-\frac{B}{4}\right)\right)+U\left(y-\frac{5 B}{8}\right)\left(-\frac{8}{B}\left(y-\frac{6 B}{8}\right)-1\right)+U\left(y-\frac{6 B}{8}\right)\left(\frac{8}{B} \times\left(y-\frac{6 B}{8}\right)\right)\right)^{2} \cos (\lambda y) d y}{\sinh (\lambda L)}$
$=\frac{2 h_{1}^{2}\left(128 \sin \left(\frac{n \pi}{4}\right)-128 \sin \left(\frac{3 n \pi}{8}\right)+128 \sin \left(\frac{5 n \pi}{8}\right)-128 \sin \left(\frac{3 n \pi}{4}\right)+16 n \pi \cos \left(\frac{3 n \pi}{8}\right)+16 n \pi \cos \left(\frac{5 n \pi}{8}\right)\right)}{a^{2} n^{3} \pi^{3} \sinh \left(\frac{n \pi L}{B}\right)}$

## Tunnel distribution

Hydraulic head contours (Fig. 9) were calculated by Eqs. (15) and (22) for $n=1 \sim 20$. As depicted in Fig. 9, the trapezoidal distribution of the hydraulic head on the rioht houndary creates a hend shane in the $h$ contours

In this case, the outlet conduit has a tunnel shape. Mathematically, the right boundary condition is rewritten as:
$h(L, y)=\frac{h_{1}}{a} \times\left(U\left(y-\frac{B}{3}\right) \times\left(-\sin \left(\frac{\pi y}{B / 3}\right)+1\right)+U\left(y-\frac{2 B}{3}\right) \times\left(\sin \left(\frac{\pi y}{B / 3}\right)-1\right)\right)$

The right boundary condition (Eq. 23) with $a=6$, $h_{l}=300 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ was shown in Fig. 10. This figure presents the tunnel distribution of the hydraulic head for $\frac{B}{3} \leq y \leq \frac{2 B}{3}$. For the tunnel distribution
case, mathematical forms of $C_{3}^{*}$ and $C_{5}^{*}$ are identical to Eq. (14) and $C_{2}^{*}$ and $C_{4}^{*}$ are expressed as:

Fig. 10 Right boundary condition (Eq. 23) for the tunnel distribution case


$$
\begin{align*}
C_{2}^{*} & =\frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(U\left(y-\frac{B}{3}\right) \times\left(-\sin \left(\frac{\pi y}{B / 3}\right)+1\right)+U\left(y-\frac{2 B}{3}\right) \times\left(\sin \left(\frac{\pi y}{B / 3}\right)-1\right)\right)^{2} d y-\frac{C_{3}^{*}}{L} \\
& =\frac{\frac{1}{6} \frac{h_{1}^{2}(3 \pi+8)}{a^{2} \pi}-h_{1}^{2}}{L} \tag{24-a}
\end{align*}
$$

$$
\begin{align*}
C_{4}^{*} & =\frac{\frac{2}{B} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} *\left(U\left(y-\frac{B}{3}\right) \times\left(-\sin \left(\frac{\pi y}{B / 3}\right)+1\right)+U\left(y-\frac{2 B}{3}\right) \times\left(\sin \left(\frac{\pi y}{B / 3}\right)-1\right)\right)^{2} \cos (\lambda y) d y}{\sinh (\lambda L)}  \tag{24-b}\\
& =\frac{-2 h_{1}^{2}}{a^{2} n \pi\left(n^{4}-45 n^{2}+324\right) \sinh \left(\frac{n \pi L}{B}\right)} \times\left(486 \sin \left(\frac{n \pi}{3}\right)-486 \sin \left(\frac{2 n \pi}{3}\right)-n^{4} \sin \left(\frac{n \pi}{3}\right)-63 n^{2} \sin \left(\frac{n \pi}{3}\right)+6 n^{3} \cos \left(\frac{n \pi}{3}\right)-n^{4} \sin \left(\frac{2 n \pi}{3}\right)\right. \\
& \left.+63 n^{2} \sin \left(\frac{2 n \pi}{3}\right)+6 n^{3} \cos \left(\frac{2 n \pi}{3}\right)-216 n \cos \left(\frac{n \pi}{3}\right)-216 n \cos \left(\frac{2 n \pi}{3}\right)\right)
\end{align*}
$$

In order to confirm convergence of summations in Eq. (15) with new coefficients, hydraulic head at different locations is calculated by summations truncation with different values of $n$ (Table 2). The dam parameters utilized for the problem are similar to the rectangular distribution case. To illustrate the application of the
derived equations, hydraulic head values from an explicit scheme FDM solution (Eq. 2) is compared to the analytical solution (Eq. 15) for different values of $x$ and $y$ (Table 2, columns 7). As presented in Table 2, the values of relative errors ( $R E s$ ) are less than $1 \%$. Based on Eqs. (15) and (24), hydraulic head contours were

Table 2 Hydraulic head values resulted from analytical approach (for different summation truncations) and FDM for tunnel distribution case

| $x(\mathrm{~cm})$ | $y(\mathrm{~cm})$ | $h_{\text {analytical }}$ |  | $h_{F D M}$ | $R E(\%)$ |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  |  | $n=1 \sim 5$ | $n=1 \sim 10$ | $n=1 \sim 15$ | $n=1 \sim 20$ | $\Delta x=5 \mathrm{~cm}$ and $\Delta y=5 \mathrm{~cm}$ |  |
| 95 | 250 | 105.57 | 115.93 | 115.68 | 115.60 | 116.69 | 0.94 |
| 90 | 250 | 128.64 | 131.27 | 131.52 | 131.57 | 132.70 | 0.85 |
| 80 | 250 | 156.21 | 157.91 | 157.72 | 157.78 | 158.92 | 0.72 |
| 70 | 250 | 179.42 | 180.53 | 180.72 | 180.78 | 181.56 | 0.43 |
| 60 | 250 | 201.17 | 201.28 | 201.49 | 201.57 | 202.64 | 0.53 |
| 20 | 250 | 270.24 | 270.39 | 270.39 | 270.39 | 270.87 | 0.17 |
| 90 | 100 | 90.12 | 95.11 | 94.95 | 95.08 | 95.24 | 0.16 |
| 80 | 350 | 137.24 | 135.61 | 135.65 | 135.68 | 135.83 | 0.11 |
| 70 | 400 | 163.25 | 164.57 | 164.60 | 164.61 | 164.86 | 0.15 |
| 30 | 400 | 250.99 | 251.17 | 251.17 | 251.17 | 251.31 | 0.05 |

Fig. 11 Hydraulic head contours for the tunnel distribution case

depicted in Fig. 11. As seen in this figure, the tunnel distribution of the hydraulic head on the right boundary of the dam plane creates a bell shape contour. Figure 11 shows the distribution of the hydraulic head that remains at $h=300 \mathrm{~cm}$ on the left $(x=0)$ and $h=$ Eq. (23) on the right boundaries $(x=100 \mathrm{~cm})$ of the dam plane.

## Piecewise rectangular distribution

Piecewise rectangular distribution can also be justified similar to the trapezoidal and tunnel distributions cases where there are several rectangular outlets instead of a single outlet. The bottom, top, and left boundary conditions for the piecewise rectangular distribution case are identical to the previous cases. The mathematical form for the right boundary condition is written as:
$h(L, y)=\frac{h_{1}}{a} \times\left(U(y)-U\left(y-\frac{B}{14}\right)+U\left(y-\frac{2 B}{14}\right)-U\left(y-\frac{3 B}{14}\right)+U\left(y-\frac{4 B}{14}\right)-U\left(y-\frac{5 B}{14}\right)+U\left(y-\frac{6 B}{14}\right)\right)$

The right boundary condition (Eq. 25) with $a=3$, $h_{l}=300 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ was depicted in Fig. 12. As depicted in this figure, the boundary condition presents the constant hydraulic head for the ranges of $0 \leq y \leq \frac{B}{14}, \frac{2 B}{14} \leq y \leq \frac{3 B}{14}, \frac{4 B}{14} \leq y \leq \frac{5 B}{14}$ and $\frac{6 B}{14} \leq y \leq \frac{7 B}{14}$. The
answer for $h(x, y)$ is similar to Eq. (15) but with different coefficients. Mathematical forms of coefficients $C_{3}^{*}$ and $C_{5}^{*}$ are identical to Eq. (14); however, coefficients $C_{2}^{*}$ and $C_{4}^{*}$ are expressed as:
$C_{2}^{*}=\frac{1}{B L} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(U(y)-U\left(y-\frac{B}{14}\right)+U\left(y-\frac{2 B}{14}\right)-U\left(y-\frac{3 B}{14}\right)+U\left(y-\frac{4 B}{14}\right)-U\left(y-\frac{5 B}{14}\right)+U\left(y-\frac{6 B}{14}\right)\right)^{2} d y-\frac{C_{3}^{*}}{L}=\frac{\frac{11}{14} \frac{h_{1}^{2}}{a^{2}}-h_{1}^{2}}{L}$
(26-a)

Fig. 12 Right boundary condition (Eq. 23) for the piecewise rectangular distribution case

$C_{4}^{*}$

$$
\begin{align*}
& =\frac{\frac{2}{B} \int_{0}^{B}\left(\frac{h_{1}}{a}\right)^{2} \times\left(U(y)-U\left(y-\frac{B}{14}\right)+U\left(y-\frac{2 B}{14}\right)-U\left(y-\frac{3 B}{14}\right)+U\left(y-\frac{4 B}{14}\right)-U\left(y-\frac{5 B}{14}\right)+U\left(y-\frac{6 B}{14}\right)\right)^{2} \cos (\lambda y) d y}{\sinh (\lambda L)}  \tag{26-b}\\
& =\frac{2 h_{1}^{2}\left(\sin \left(\frac{n \pi}{14}\right)-\sin \left(\frac{n \pi}{7}\right)+\sin \left(\frac{3 n \pi}{14}\right)-\sin \left(\frac{2 n \pi}{7}\right)+\sin \left(\frac{5 n \pi}{14}\right)-\sin \left(\frac{3 n \pi}{7}\right)\right)}{a^{2} n \pi \sinh \left(\frac{n \pi L}{B}\right)}
\end{align*}
$$

Fig. 13 Hydraulic head contours for the piecewise rectangular distribution case



Fig. 14 River bank protection wall as a case for the 3D water seepage analysis

Based on Eqs. (15) and (24), hydraulic head contours were illustrated in Fig. 13. As seen in this figure, the piecewise rectangular distribution of the hydraulic head on the right boundary creates a wrinkle shape in the $h$ contours.

## Analytical solution for 3D water seepage

Seepage analysis in the river bank (Javdanian and Jafarian 2018) protection walls is an example of the 3D water seepage (Fig. 14). As shown in Fig. 14, on the left side and upstream of the wall is collected water that can be

Downstream side


Fig. 15 Solution domain for 3D water seepage analysis with a linear distribution of the hydraulic head at downstream
accumulated because of raining water, sewage water, dirty water, and so forth. The collected water creates a constant hydraulic head at the upstream of the wall. The river has a natural flow on the downstream of the wall. As shown in Fig. 14, the river flow due to the hydraulic gradient creates a linear hydraulic head on the downstream of the wall. Therefore, a constant hydraulic head at the upstream and a linear distribution of the hydraulic head at the downstream of the dam (Fig. 15) were applied. Due to the 3D domain and high hydraulic gradient at the downstream of the dam, there are both vertical and horizontal seepage. As a result, the Dupuit assumptions are not suitable for this case. The governing equation for 3D steady-state water seepage in the isotropic homogeneous dam without Dupuit assumptions is as follows (see Eqs. 1 and 2):
$\frac{\partial^{2} h^{2}}{\partial x^{2}}+\frac{\partial^{2} h^{2}}{\partial y^{2}}+\frac{\partial^{2} h^{2}}{\partial z^{2}}=0$
where h is the hydraulic head $(1 / L)$. In this case, the no-flow boundary condition is applied to three sides of the dam. Boundary conditions for this case are mathematically expressed as:
$\frac{\partial h}{\partial z}(x, y, 0)=0 \quad h\left(x, y, h_{1}\right)=0$
$\frac{\partial h}{\partial y}(x, 0, z)=0 \quad \frac{\partial h}{\partial y}(x, B, z)=0$
$h(0, y, z)=h_{1}$
$h(L, y, z)=\left[U(z)-U\left(z-\left(\frac{h_{2}-h_{3}}{B} y+h_{3}\right)\right)\right] \times\left[\frac{h_{2}-h_{3}}{B} y+h_{3}\right]$
where $U$ is the unit step function. Also, $h_{1}, h_{2}$, and $h_{3}$ were shown in Fig. 15. The downstream boundary condition for all points below the water table has a linear distribution of the hydraulic head, and for all points above the water table provides zero hydraulic head. Also, this boundary condition provides a constant hydraulic head value for all points located on any vertical line that is perpendicular to the lower edge of the dam at the downstream side.

To find an analytical solution for Eq. (27) with the aforementioned boundary conditions, the variable change
as $h^{2}(x, y, z)=w(x, y, z)$ is used, and then the PDE and the boundary conditions for $w(x, y, z)$ is written as:
$\frac{\partial^{2} w}{\partial x^{2}}+\frac{\partial^{2} w}{\partial y^{2}}+\frac{\partial^{2} w}{\partial z^{2}}=0$
$\frac{\partial w}{\partial z}(x, y, 0)=0 \quad w\left(x, y, h_{1}\right)=0$
$\frac{\partial w}{\partial y}(x, 0, z)=0 \quad \frac{\partial w}{\partial y}(x, B, z)=0$
$w(0, y, z)=h_{1}{ }^{2}$
$w(L, y, z)=\left(\left[U(z)-U\left(z-\left(\frac{h_{2}-h_{3}}{B} y+h_{3}\right)\right)\right] \times\left[\frac{h_{2}-h_{3}}{B} y+h_{3}\right]\right)^{2}$

Using separation of variables for $w(x, y, z)$ :
$w(x, y, z)=X(x) Y(y) Z(z)$
Substituting Eq. (31) into Eq. (29) gives:
$\frac{Z^{\prime \prime}}{Z}=-\left(\frac{X^{\prime \prime}}{X}+\frac{Y^{\prime \prime}}{Y}\right)=\mu$
If $\mu \geq 0$, a trivial solution for $Z(z)$ in Eq. (32) would be obtained. If $\mu<0$, assuming $\mu=-\lambda^{2}, \lambda>0$, then by using the boundary conditions of Eq. (30-a), $Z(z)$ in Eq. (32) is written as:
$Z_{n}=B_{2} \cos (\lambda z), \quad \lambda=\frac{(2 n-1) \pi}{2 h_{1}}, n=1,2,3, \ldots, \infty$
where $B_{2}$ is a constant.
Substituting $-\lambda^{2}$ into Eq. (32) for $\frac{Z^{\prime \prime}}{Z}$ yields:
$\frac{Y^{\prime \prime}}{Y}=\lambda^{2}-\frac{X^{\prime \prime}}{X}=\rho$
where $\rho$ is an arbitrary constant.
If $\rho=0$ and $\rho<0$, say $\rho=-\beta^{2}$ and $\beta>0$, then by considering the boundary conditions of Eq. (30-b), $Y(y)$ in Eq. (34) has two answer as follows:
$Y_{1}=D$
$Y_{n}=B_{3} \cos (\beta y), \quad \beta=\frac{m \cdot \pi}{B}, m=1,2,3, \ldots, \infty$
where $D$ and $B_{3}$ are constants. For $\rho>0$, a trivial solution is obtained for $Y(y)$. Substituting $\rho=0$ and $\rho=-\beta^{2}$ in Eq. (34) for $\frac{Y^{\prime} \prime}{Y}$, two equations are obtained:
$\lambda^{2}-\frac{X^{\prime \prime}}{X}=0$
$\lambda^{2}-\frac{X^{\prime \prime}}{X}=-\beta^{2}$
The solution of Eq. (37-a) and Eq. (37-b) gives:
$X_{n}=A_{5} \cosh (\lambda x)+A_{6} \sinh (\lambda x)$
$X_{n m}=A_{4} \cosh (\tau x)+B_{4} \sinh (\tau x), \quad \tau=\sqrt{\left(\beta^{2}+\lambda^{2}\right)}$

By substituting Eqs. (33), (36), and (38-b) into Eq. (31), the second answer for $w(x, y, z)$ is written as:

$$
\begin{align*}
w_{2}(x, y, z) & =X_{n m}(x) Y_{n}(y) Z_{n}(z) \\
& =\sum_{n=1}^{\infty} \sum_{m=1}^{\infty}\left(A_{4}{ }^{*} \cosh (\tau x)+\mathrm{B}_{4}^{*} \sinh (\tau x)\right) \cos (\beta y) \cos (\lambda z) \tag{40}
\end{align*}
$$

where $A_{4}{ }^{*}$ is $A_{4} B_{2} B_{3}$ and $\mathrm{B}_{4}{ }^{*}$ is $B_{4} B_{2} B_{3}$.
Finally, the overall answer for $w(x, y, z)$ is the combination of Eq. (39) and Eq. (40):

$$
\begin{align*}
w(x, y, z) & =w_{1}(x, y, z)+w_{2}(x, y, z)  \tag{41}\\
& =\sum_{n=1}^{\infty} \sum_{m=1}^{\infty}\left(A_{4}{ }^{*} \cosh (\tau x)+\mathrm{B}_{4}^{*} \sinh (\tau x)\right) \cos (\beta y) \cos (\lambda z) \\
& +\sum_{n=1}^{\infty}\left[A_{5}{ }^{*} \cosh (\lambda x)+A_{6}{ }^{*} \sinh (\lambda x)\right] \cos (\lambda z)
\end{align*}
$$

Substituting the boundary condition of Eqs. (30-c) and (30-d) into Eq. (41) and using Fourier series properties gives:
where $A_{4}, A_{5}, A_{6}$, and $B_{4}$ are constants.
By substituting Eqs. (33), (35), and (38-a) into Eq. (31), the first answer for $w(x, y, z)$ yields:

$$
\begin{align*}
& w_{1}(x, y, z)=X_{n}(x) Y_{1}(y) Z_{n}(z)  \tag{42-b}\\
& \quad=\sum_{n=1}^{\infty}\left[A_{5}{ }^{*} \cosh (\lambda x)+A_{6}{ }^{*} \sinh (\lambda x)\right] \cos (\lambda z) \tag{39}
\end{align*}
$$

$$
A_{5}^{*}=\frac{2}{h_{1} B} \int_{0}^{B} \int_{0}^{h_{1}} h_{1}^{2} \cos (\lambda z) d z d y=\frac{4 h_{1}^{2}(-1)^{n+1}}{(2 n-1) \pi}
$$

where $A_{5}{ }^{*}$ is $A_{5} D B_{2}$ and $A_{6}{ }^{*}$ is $A_{6} D B_{2}$.
$B_{4}{ }^{*}=\frac{1}{\sinh (\tau L)} * \frac{4}{h_{1} B} \int_{0}^{B} \int_{0}^{h_{1}}\left(\left[U(z)-U\left(z-\left(\frac{h_{2}-h_{3}}{B} y+h_{3}\right)\right)\right] \times\left[\frac{h_{2}-h_{3}}{B} y+h_{3}\right]\right)^{2} \cos (\beta y) \cos (\lambda z) d z d y$
$A_{6}{ }^{*}=\frac{\frac{2}{h_{1} B} \int_{0}^{B} \int_{0}^{h_{1}}\left(\left[U(z)-U\left(z^{-}\left(\frac{h_{2}-h_{3}}{B} y+h_{3}\right)\right)\right] \times\left[\frac{h_{2}-h_{3}}{B} y+h_{3}\right]\right)^{2} \cos (\lambda z) d z d y-A_{5}{ }^{*} \cosh (\lambda L)}{\sinh (\lambda L)}$

Table 3 Hydraulic head values resulted from analytical approach (for different summation truncations) and FDM for 3D water seepage at the vertical section $y=250 \mathrm{~cm}$

| $x(\mathrm{~cm})$ | $z(\mathrm{~cm})$ | $h_{\text {analytical }}$ |  |  |  | $h_{F D M}$ | $R E$ (\%) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $\begin{aligned} & n=m= \\ & 1 \sim 5 \end{aligned}$ | $\begin{aligned} & n=m= \\ & 1 \sim 10 \end{aligned}$ | $\begin{aligned} & n=m= \\ & 1 \sim 15 \end{aligned}$ | $\begin{aligned} & n=m= \\ & 1 \sim 20 \end{aligned}$ | $\Delta x=\Delta y=\Delta z=10 \mathrm{~cm}$ |  |
| 30 | 100 | 260.92 | 261.58 | 262.05 | 262.05 | 265.34 | 1.25 |
| 50 | 150 | 223.14 | 223.81 | 224.37 | 224.37 | 227.02 | 1.18 |
| 60 | 250 | 159.48 | 160.02 | 160.69 | 160.69 | 163.21 | 1.56 |
| 80 | 200 | 131.56 | 132.21 | 132.83 | 132.83 | 134.62 | 1.34 |
| 90 | 250 | 76.31 | 77.14 | 77.73 | 77.73 | 79.01 | 1.64 |
| 20 | 150 | 270.68 | 271.23 | 271.84 | 271.84 | 274.35 | 0.92 |
| 10 | 270 | 263.09 | 263.76 | 264.20 | 264.20 | 267.41 | 1.21 |
| 40 | 180 | 233.25 | 233.87 | 234.47 | 234.47 | 237.22 | 1.17 |
| 50 | 250 | 182.08 | 182.71 | 183.20 | 183.20 | 186.31 | 1.69 |
| 90 | 150 | 127.33 | 127.83 | 128.37 | 128.37 | 130.85 | 1.93 |



Fig. 16 Hydraulic head contours for 3D water seepage on the vertical sections at: $\mathbf{a} y=20 \mathrm{~cm}, \mathbf{b} y=250 \mathrm{~cm}, \mathbf{c} y=480 \mathrm{~cm}$

Some of the terms in the answer of the integral Eq. (42-c) with $h_{1}=300 \mathrm{~cm}, h_{2}=200 \mathrm{~cm}, h_{3}=$ 100 cm , and $B=500 \mathrm{~cm}$ are as follows:

$$
\begin{align*}
& \int_{0}^{500} \int_{0}^{300}\left(\left[U(z)-U\left(z^{-}\left(\frac{200-100}{500} y+100\right)\right)\right] \times\left[\frac{200-100}{500} y+100\right]\right)^{2} \cos \left(\frac{m \pi}{500} y\right) \cos \left(\frac{(2 n-1) \pi}{2 * 300} z\right)  \tag{43}\\
& d z d y=9 * 10^{9} *\left(4(-1)^{1+m} \cos \left(\frac{2}{3} \pi n\right) \pi^{2}+24(-1)^{1+m} \sin \left(\frac{2}{3} \pi n\right) \pi+432(-1)^{1+m} n \cos \left(\frac{2}{3} \pi n\right)\right. \\
&+576(-1)^{1+m} n^{3} \cos \left(\frac{2}{3} \pi n\right)+768(-1)^{1+m} n^{3} \cos \left(\frac{2}{3} \pi n\right) \sqrt{3} \pi+320(-1)^{1+m} n^{4} \sin \left(\frac{2}{3} \pi n\right) \sqrt{3} \pi^{2} \\
&+192(-1)^{1+m} n \cos \left(\frac{2}{3} \pi n\right) \sqrt{3} \pi+2304(-1)^{1+m} n^{3} \pi^{2} m^{2} \cos \left(\frac{2}{3} \pi n\right)+1728(-1)^{1+m} n \pi^{2} m^{2} \cos \left(\frac{2}{3} \pi n\right) \\
&\left.+5184(-1)^{1+m} \pi^{2} m^{4} \sin \left(\frac{2}{3} \pi n\right) \sqrt{3}+\ldots\right) \\
& /\left(\pi^{4}\left(1+448 n^{6}-672 n^{5}+560 n^{4}-46656 m^{6}-280 n^{3}+3888 m^{4}+84 n^{2}-108 m^{2}+3456 n^{5} m^{2}+\ldots\right)\right)
\end{align*}
$$

Also, some of the terms in the answer of the integral Eq. (42-d) with $h_{1}=300 \mathrm{~cm}, h_{2}=200 \mathrm{~cm}, h_{3}=100 \mathrm{~cm}$, and $B=500 \mathrm{~cm}$ are written as:

$$
\begin{align*}
& \int_{0}^{500} \int_{0}^{300}\left(\left[U(z)-U\left(z^{-}\left(\frac{200-100}{500} y+100\right)\right)\right] \times\left[\frac{200-100}{500} y+100\right]\right)^{2} \cos \left(\frac{(2 n-1) \pi}{2 * 300} z\right)  \tag{44}\\
& d z d y= \frac{-18 * 10^{9}}{(2 n-1)^{4} \pi^{4}} \times\left(-72 \cos \left(\frac{1}{3}(2 n-1) \pi\right)+72 \cos \left(\frac{1}{6}(2 n-1) \pi\right)+4 \cos \left(\frac{1}{3}(2 n-1) \pi\right) \pi^{2}\right. \\
&\left.\quad+16 \cos \left(\frac{1}{3}(2 n-1) \pi\right) \pi^{2} n^{2}-16 \cos \left(\frac{1}{3}(2 n-1) \pi\right) \pi^{2} n+24 \pi \sin \left(\frac{1}{3}(2 n-1) \pi\right)-48 \pi n \sin \left(\frac{1}{3}(2 n-1) \pi\right)+\ldots\right)
\end{align*}
$$

Substituting Eq. (41) into $h^{2}(x, y, z)=w(x, y, z)$ gives:

$$
\begin{aligned}
& h^{2}(x, y, z)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty}\left(A_{4}{ }^{*} \cosh (\tau x)+\mathrm{B}_{4}^{*} \sinh (\tau x)\right) \cos (\beta y) \\
& \quad \cos (\lambda z)+\sum_{n=1}^{\infty}\left[A_{5}{ }^{*} \cosh (\lambda x)+A_{6}{ }^{*} \sinh (\lambda x)\right] \\
& \\
& \cos (\lambda z)
\end{aligned}
$$

In order to test convergence of summations in Eq. (45), hydraulic head at a vertical section ( $y=$ 250 cm ) is calculated by summations truncations with different amounts of $n$ and $m$ (Table 3). Table 3 is presented for the parameters $h_{1}=300 \mathrm{~cm}, h_{2}=200 \mathrm{~cm}$, $h_{3}=100 \mathrm{~cm}, B=500 \mathrm{~cm}$, and $L=100 \mathrm{~cm}$.

To illustrate application of the derived equations, hydraulic head values from an explicit scheme FDM solution (Eq. 27) is compared to the analytical solution (Eq. 45) for different values of $x$ and $z$ (Table 3, columns 7). The eighth column depicts relative error $(R E)$ based on column 6 (incorporating $>400$ summation terms) and column 7 (FDM for $\Delta x=\Delta y=\Delta z=10 \mathrm{~cm}$ ). As presented in Table 3, the values of $R E$ are less than $2 \%$.

Based on Eq. (45), hydraulic head contours were depicted in Figs. 16(a-c) at the vertical sections of $y=$ 20,250 , and 480 cm , respectively. As seen in these figures, by increasing the $y$ value, hydraulic head increases in all parts of the vertical sections ( $x-z$ plane) of the dam (corresponding to Fig. 15 and Eq. 28-d).

Figure 16 illustrates distribution of the hydraulic head that remains at $h=300 \mathrm{~cm}$ on the left $(x=0)$ and $h=$ Eq. ( $28-\mathrm{d}$ ) on the right ( $x=100 \mathrm{~cm}$ ) boundaries in the vertical sections ( $x-z$ plane) of the dam.

## Summary and conclusions

Innovative analytical solutions for two- and threedimensional steady-state water seepage through dams with nonsymmetric boundary conditions were derived using separation variables and Fourier series expansion techniques. The nonsymmetric boundary conditions were considered by different unit step functions on a part and/or parts of the boundary of the dam plane. Six two-dimensional cases were studied, where a constant hydraulic head was applied at the left boundary and various hydraulic head distributions including rectangular, ramp, triangular, trapezoidal, tunnel, and piecewise rectangular distributions were applied at the right boundary of the dam plane. Subsequently, a threedimensional case was investigated so that a constant hydraulic head was applied at the upstream and a linearly distributed hydraulic head was applied at the downstream of the dam. All the boundary conditions and partial differential equations (PDEs) were satisfied by the derived analytical solutions. Subsequently, the performance of the developed analytical methods was assessed in comparison with numerical modeling. The results illustrate reasonable accuracy of the governing equations to assess seepage through media with nonsymmetric boundary conditions. The developed analytical solutions can be used as a reference result to evaluate accuracy and verify the numerical algorithms with similar nonsymmetric boundary conditions.

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